

# Contents

<b>1</b>	<b>Pre-mid</b>	<b>3</b>
<b>2</b>	<b>Week 7</b>	<b>4</b>
2.1	Sampling Theorem . . . . .	4
<b>3</b>	<b>Week 8</b>	<b>5</b>
3.1	Wavelets and multiresolution analysis . . . . .	5
3.1.1	Introduction . . . . .	5
<b>4</b>	<b>Week 9</b>	<b>6</b>
4.1	Intro cont'd . . . . .	6
4.1.1	Nesting relation . . . . .	6
4.1.2	Nesting relation $V_0 \subset V_1$ . . . . .	6
4.2	wavalets . . . . .	7
4.2.1	Summary of major recent results . . . . .	7
4.2.2	Attention! Important summary . . . . .	8
4.2.3	Higher-order nestings $V_j \subset V_{j+1}$ . . . . .	8
4.2.4	Connection to the (continuous) wavelet transform introduced earlier . . . . .	8
4.2.5	Special case: Haar wavelet expansions of functions on a finite interval . . . . .	9
4.2.6	Analysis and synthesis alg for wavelet expansions . . . . .	9

<b>5</b>	<b>Week 10</b>	<b>11</b>
5.1	Analysis and synthesis alg for wavelet expansions cont'd . . . . .	11
5.1.1	General Resolutions . . . . .	11
5.2	Multi-re analysis: A general treatment . . . . .	11
5.2.1	Wavelet spaces . . . . .	12
<b>6</b>	<b>Week 11</b>	<b>14</b>
6.1	MRA: general (cont'd) . . . . .	14
6.1.1	Wavelet spaces (cont'd) . . . . .	14
6.1.2	Synthesis and Analysis Alg for MRAs . . . . .	14
6.1.3	Wavelets with compact support . . . . .	15
6.1.4	Relating the support of $\phi(x)$ to nonzero $h_k$ coeff . . . . .	18
<b>7</b>	<b>Week 12</b>	<b>20</b>
7.1	MRA: general treatment (cont'd) . . . . .	20
7.1.1	$\psi(x)$ its support, and vanishing moments . . . . .	20
7.1.2	Vanishing moments and the approximation of functions . . . . .	22
7.2	MRA and Fourier transform . . . . .	22
7.2.1	Fourier transforms and vanishing moments of wavelets . . . . .	22

# ♠ | Pre-mid

Review the proofs for fourier relations...



## Week 7

### 2.1 Sampling Theorem

$\Omega$  is the angular velocity.

A function  $f(t)$  is said to be bandlimited, or  $\Omega$ -bandlimited, if there exists an  $\Omega > 0$  such that

$$F(\omega) = 0 \quad \text{for } |\omega| > \Omega$$

**Nyquist freq:**  $\nu = \frac{\Omega}{2\pi}$

**The Whittaker-Shannon Sampling Theorem** :  $f(t)$   $\Omega$ -bandlimited. Then  $f = \mathcal{F}^{-1}F$  is completely determined at any  $t \in \mathbb{R}$  by its values at  $t_k = \frac{k\pi}{\Omega}$ ,  $k = 0, \pm 1, \pm 2, \dots$ , as follows

$$f(t) = \sum_{k=-\infty}^{\infty} f\left(\frac{k\pi}{\Omega}\right) \frac{\sin(\Omega t - k\pi)}{\Omega t - k\pi} = \sum_{k=-\infty}^{\infty} f\left(\frac{k\pi}{\Omega}\right) \operatorname{sinc}\left(\frac{\Omega t}{\pi} - k\right)$$

**The sampling freq is twice the bandwidth freq  $\Omega$**

effects of undersampling



## Week 8

one consequence of uncertainty principle: A function and its Fourier transform cannot both have finite support.

windowed FT, local freq of a signal

### 3.1 Wavelets and multiresolution analysis

#### 3.1.1 Introduction

Detail or wavelet function

$$\psi(t) = \begin{cases} 1 & 0 \leq t < 1/2 \\ -1 & 1/2 \leq t < 1 \end{cases}$$

Let us introduce the following space of functions:

$$V_0 = \{f \in L^2(\mathbb{R}) : f(t) \text{ is constant over the interval } [k, k+1), \quad \forall k \in \mathbb{Z}\}$$

and  $\phi_{0k}(t) = I_{[k, k+1)}(t)$  spans  $V_0$ .

$$\phi_{0k}(t) = \phi(t - k)$$

We notice that the  $\phi_{1k}$  are translated copies of  $\phi_{10}$ . But what is more important is that they also are **dilated and translated** copies of the scaling function  $\phi(t)$ :

$$\phi_{1k} = \sqrt{2}\phi(2t - k)$$

they span  $V_1$ .



## Week 9

### 4.1 Intro cont'd

Go backwards:

$$V_{-1} = \dots \text{ constant over } [2k, 2k+1) = \text{span} \left\{ \frac{1}{\sqrt{2}} \phi \left( \frac{t}{2} - k \right), \forall k \in \mathbb{Z} \right\} \cap L^2(\mathbb{R})$$

In general, for  $J \in \mathbb{Z}$ .

$$V_J = \{ \dots \text{over } \left[ \frac{k}{2^J}, \frac{k+1}{2^J} \right), \forall k \in \mathbb{Z} \}$$

and it is spanned by  $2^{J/2} \phi(2^J t - k)$ .

#### 4.1.1 Nesting relation

$$\dots \subset V_{j-1} \subset V_j \subset V_{j+1} \subset \dots$$

#### 4.1.2 Nesting relation $V_0 \subset V_1$

$$\phi(t) = \phi(2t) + \phi(2t-1), \quad t \in [0, 1]$$

More mathematical:

$$\phi(t) \sum_{k \in \mathbb{Z}} h_k \phi_{1k}(t) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2t - k)$$

latter one is called multiresolution analysis or refinement equation that is satisfied by a scaling function. It is sometimes simply called scaling equation. The nonzero coeff  $h_k$  are known as scaling coeff. For Haar system:

$$h_k = \begin{cases} \frac{1}{\sqrt{2}} & k = 0, 1 \\ 0 & \text{otherwise} \end{cases}$$

## 4.2 wavalets

Using  $V_0 \subset V_1$ , we define  $W_0 = V_0^\perp$ .

So  $V_1 = V_0 \oplus W_0$ .

The function  $\phi(t)$  is known as the **mother wavelet** of the Haar wavelet system, or simply the **Haar mother wavelet**.

### 4.2.1 Summary of major recent results

1. multiresolution analysis:

$$\phi(t) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2t - k)$$

non-zero coeffs characterize a particular multiresolution analysis.

2. For each  $j \in \mathbb{Z}$ , the infinite set of functions

$$\phi_{jk}(t) = 2^{j/2} \phi(2^j t - k), \quad k \in \mathbb{Z}$$

orthonormal basis for  $V_j$ .

3. nested

4.  $V_1 = V_0 \oplus W_0$

5. **Haar wavelet function**  $\psi(t)$  is a function in  $V_1$  that is orthogonal to  $\phi(t)$

$$\psi(t) = \phi(2t) - \phi(2t - 1)$$

More generally,

$$\psi(t) = \sum_{k \in \mathbb{Z}} g_k \sqrt{2} \phi(2t - k)$$

6. The integer translates of  $\psi(t)$  form an orthonormal basis of  $W_0$ .

A consequence  $V_1 = V_0 \oplus W_0$ , we need two sets of functions:  $\phi_{0k}$  and  $\psi_{0k}$ :

$$u(t) = \sum_{k \in \mathbb{Z}} a_k \phi_{0k}(t) + \sum_{k \in \mathbb{Z}} b_k \psi_{0k}(t)$$

### 4.2.2 Attention! Important summary

The function  $f_1 \in V_1$  which is the best approximation of  $f$  in the space  $V_1$  will admit the following expansion

$$f_1(t) = \sum_{k \in \mathbb{Z}} a_k \phi_{0k}(t) + \sum_{k \in \mathbb{Z}} b_k \psi_{0k}(t)$$

where

$$a_k = \langle f_1, \phi_{0k} \rangle, \quad b_k = \langle f_1, \psi_{0k} \rangle, \quad k \in \mathbb{Z}$$

Also  $f_1 \in V_1$ :

$$f_1(t) = \sum_{k \in \mathbb{Z}} c_k \phi_{1k}(t)$$

where

$$c_k = \langle f_1, \phi_{1k} \rangle$$

### 4.2.3 Higher-order nestings $V_j \subset V_{j+1}$

$$V_{j+1} = V_j \oplus W_j$$

Then

$$\lim_{J \rightarrow \infty} V_J = L^2(\mathbb{R})$$

This means that

$$L^2(\mathbb{R}) = V_0 \oplus \left[ \bigoplus_{j=0}^{\infty} W_j \right]$$

If we decompose  $V_0$  more, then

$$L^2(\mathbb{R}) = \bigoplus_{j=0}^{\infty} W_j$$

Consequently, any function  $f \in L^2(\mathbb{R})$  admits a unique expansion of the form

$$f(t) = \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}} b_{jk} \psi_{jk}(t)$$

### 4.2.4 Connection to the (continuous) wavelet transform introduced earlier

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#### 4.2.5 Special case: Haar wavelet expansions of functions on a finite interval

$$f(t) = f_0(t) + \sum_{j=1}^{\infty} w_j(t)$$

where  $f_0(t) = a_{00}\phi_{00}(t)$ . Detail functions,  $w_j \in W_j$  are defined:

$$w_j(t) = \sum_{k=0}^{2^j-1} b_{jk}\psi_{jk}(t), \quad j = 0, 1, 2, \dots$$

In **practical** situations, we deal with finite-dimensional.

$$f_j(t) = f_0(t) + \sum_{i=0}^{j-1} w_i(t) = a_{00}\phi_{00}(t) + \sum_{i=0}^{j-1} \sum_{k=0}^{2^i-1} b_{ik}\psi_{ik}(t)$$

**Summary** <sup>1</sup>  $a_{00}$  represents constant approx.  $f_0 \in V_0$  to  $f$ . Using  $a_{00}$  and  $b_{00}$  produces  $f_1$ . Then adding  $b_{10}$  and  $b_{11}$  produces  $f_2 \in V_2$ .

#### 4.2.6 Analysis and synthesis alg for wavelet expansions

**Analysis/decomposition alg**

$$f_1(t) = \sum_{k \in \mathbb{Z}} a_{1k} \sqrt{2} \phi(2t - k)$$

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<sup>1</sup>last paragraph of page 318

and

$$f_1(t) = \sum_{k \in \mathbb{Z}} a_{0k} \phi(t - k) + \sum_{k \in \mathbb{Z}} b_{0k} \psi(t - k)$$

Now introduce  $A_{jk} = 2^{j/2} a_{jk}$ ,  $B_{jk} = 2^{j/2} b_{jk}$

Then two eqs above become:

$$f_1(t) = \sum_{k \in \mathbb{Z}} A_{1k} \phi(2t - k)$$

and

$$f_1(t) = \sum_{k \in \mathbb{Z}} A_{0k} \phi(t - k) + \sum_{k \in \mathbb{Z}} A_{0k} \psi(t - k)$$

After some calculations, we get

$$A_{0k} = \frac{1}{2} [A_{1,2k} + A_{1,2k+1}]$$

$$B_{0k} = \frac{1}{2} [A_{1,2k} - A_{1,2k+1}]$$

Given coeff for  $f_1$ , we compute coeff of  $V_0 \oplus W_0$  decomposition of  $V_1$ .



## Week 10

### 5.1 Analysis and synthesis alg for wavelet expansions cont'd

#### 5.1.1 General Resolutions

$$V_j = V_{j-1} \oplus W_{j-1}$$

$$A_{j-1,k} = \frac{1}{2}[A_{j,2k} + A_{j,2k+1}]$$

$$B_{j-1,k} = \frac{1}{2}[A_{j,2k} - A_{j,2k+1}]$$

#### Reconstruction/synthesis

**Step  $j$**  Take  $a_{j-1,k}$  coeff, along with  $b_{jk}$ . compute  $a_{jk}$ ,  $0 \leq k \leq 2^j - 1$ .

### 5.2 Multi-re analysis: A general treatment

The collection  $\{V_j\}$  is called a multi-re analysis with scaling function  $\phi$  if the following conditions hold:

1. nesting

2. density:  $\overline{\bigcup_{j \in \mathbb{Z}} V_j} = L^2(\mathbb{R})$

This essentially states, in proper set-theoretic language, that  $\lim_{j \rightarrow \infty} V_j = L^2(\mathbb{R})$

3. separation:  $\bigcap_{j \in \mathbb{Z}} V_j = \{0\}$

$$4. \text{ scaling: } f(x) \in V_j \iff f(2x) \in V_{j+1}$$

$$f(x) \in V_j \iff f(2^{-j}x) \in V_0$$

Subspaces  $V_j$  satisfying 1-4 are known as approximation spaces

5. orthonormal basis

**Theorem** If the support of the scaling function  $\phi(x)$  is finite, then only a finite number of the coeff  $h_k$  can be nonzero.

**Proof** Suppose  $\phi(x) = 0$  outside the interval  $[-a, a]$ , where  $a > 0$  is finite. Also let  $k_1 < k_2 < \dots$  be an infinite seq of ints for which  $h_{k_i} \neq 0$ . Now suppose that  $\phi(p) \neq 0$  for some  $p \in [-a, a]$ . Then from the scaling eqs,

$$\phi(x) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2x - k)$$

it follows that there will be nonzero contributions to RHS at the points  $x_i \in \mathbb{R}$  defined by  $2x_i - k_i = p$ ,  $i = 1, 2, \dots$ , implying that the values  $\phi(x_i)$  are nonzero. But a rearrangement yields  $x_i = \frac{1}{2}(p + k_i)$ , implying that  $x_i \rightarrow \infty$  as  $i \rightarrow \infty$ . This contradicts the assumption that  $\phi(x)$  is zero outside the interval  $[-a, a]$ .  $\square$

### 5.2.1 Wavelet spaces

$$W_0 = V_0^\perp$$

$$\psi(x) = \sum_{k \in \mathbb{Z}} g_k \phi_{1k} = \sum_{k \in \mathbb{Z}} g_k \sqrt{2} \phi(2x - k)$$

and

$$\phi(x) = \sum_{k \in \mathbb{Z}} h_k \phi_{1k} = \sum_{k \in \mathbb{Z}} \sqrt{2} \phi(2x - k)$$

Inner product = 0

$$\langle \psi, \phi \rangle = \sum_{k \in \mathbb{Z}} \sum_{l \in \mathbb{Z}} g_k \bar{h}_l \langle \phi_{1k}, \phi_{1l} \rangle = \sum_{k \in \mathbb{Z}} g_k \bar{h}_k$$

It must = 0 by orthogonality.

**cheap trick:** Set  $g_k = (-1)^k \bar{h}_{1-k}$

Then

$$\begin{aligned}
 \sum_{k \in \mathbb{Z}} g_k \bar{h}_k &= \dots + g_{-1} \bar{h}_{-1} + g_0 \bar{h}_0 + g_1 \bar{h}_1 + g_2 \bar{h}_2 + \dots \\
 &= \dots - \bar{h}_2 \bar{h}_{-1} + \bar{h}_1 \bar{h}_0 - \bar{h}_0 \bar{h}_1 + \bar{h}_2 \bar{h}_{-1} + \dots \\
 &= 0
 \end{aligned}$$

**Theorem** For any  $j \in \mathbb{Z}$ , the set of functions  $\{\psi_{jk} = 2^{j/2} \psi(2^j x - k)\}$  where

$$\psi(x) = \sum_{k \in \mathbb{Z}} (-1)^k \bar{h}_{1-k} \phi(2x - k)$$

forms an orthonormal basis of  $W_j$ .



# Week 11

## 6.1 MRA: general (cont'd)

### 6.1.1 Wavelet spaces (cont'd)

From density property for MRA, we may write loosely:  $\lim_{j \rightarrow \infty} V_j = L^2(\mathbb{R})$ , so that the equation becomes

$$L^2(\mathbb{R}) = V_0 \oplus W_0 \oplus \dots$$

Then...

$$L^2(\mathbb{R}) = \bigoplus_{j=-\infty}^{\infty} W_j$$

The **doubly indexed** set of functions  $\{\psi_{jk} = 2^{j/2}\psi(2^j - k)\}$ ,  $j, k \in \mathbb{Z}$  forms an orthonormal basis in  $L^2(\mathbb{R})$ .

### 6.1.2 Synthesis and Analysis Alg for MRAs

#### Analysis

General decomposition:

$$V_j = V_{j-1} \oplus W_{j-1}$$

finer scale = coarser scale + detail.

In scaling function, replace  $x$  with  $2^{j-1}x - l$ , and  $m = 2l + k$ , we have

$$\phi(2^{j-1}x - l) = \sum_{m \in \mathbb{Z}} h_{m-2l} \sqrt{2} \phi(2^j x - m)$$

**Analysis** Express the coarser coeff  $a_{j-1,k}$ ,  $b_{j-1,k}$  in terms of finer coeff  $a_{j,k}$ .

Then calculate directly using inner product

$$a_{j-1,l} = \langle f, \phi_{j-1,l} \rangle$$

**Synthesis** express the finer coeff  $a_{j,k}$  in terms of coarser coeff  $a_{j-1,k}, b_{j-1,k}$ .

### 6.1.3 Wavelets with compact support

**very important!**

Two ways:

$$\phi(x) = \sum_{k \in \mathbb{Z}} h_k \phi_{1k}, \quad h_k = \langle \phi, \phi_{1k} \rangle$$

and

$$\phi(x) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2x - k)$$

Aforementioned theorem: finite support of  $\phi$  implies finite number of nonzero coeff  $h_k$ .

some conditions that must be satisfied by the scaling coeff  $h_k$  for  $\phi(x)$  to have compact support.

1. Finite energy (squared  $L^2$  norm)

$$\langle \phi, \phi \rangle = \sum_k |h_k|^2 = 1$$

2. Finite  $L^1$  norm.

$$\phi(x) = \sum_{k \in \mathbb{Z}} h_k \sqrt{2} \phi(2x - k) \quad ^1$$

Since in  $L^2[a, b]$ , then  $L^1[a, b]$

$$\left| \int_{\mathbb{R}} \phi(x) dx \right| \leq \int_{\mathbb{R}} |\phi(x)| dx < \infty$$

implying that the integral on the left exists. Now integrate both sides of scaling equation:

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<sup>1</sup>An important eq which I will use later

$$\int_{\mathbb{R}} \phi(x) dx = \sum_k h_k \sqrt{2} \int_{\mathbb{R}} \phi(2x - k) dx$$

For each  $k$ , we have this by let  $s = 2x - k$ ,

$$\int_{\mathbb{R}} \phi(2x - k) dx = \frac{1}{2} \int_{\mathbb{R}} \phi(s) ds$$

Then sub into prev eq,

$$\int_{\mathbb{R}} \phi(x) dx = \sum_k h_k \frac{1}{\sqrt{2}} \int_{\mathbb{R}} \phi(x) dx$$

Since non-zero integral, then we have

$$\sum_k h_k = \sqrt{2}$$

### 3. Generalized orthogonality

$$\langle \phi(x), \phi(x - p) \rangle = 2 \sum_k \sum_l h_k h_l \langle \phi(2x - k), \phi(2x - 2p - l) \rangle = \delta_{0p}$$

Setting  $k = 2p + l$ , we make the inner product inside to be  $\frac{1}{2}$  since  $\langle \phi(2x - k), \phi(2x - 2p - l) \rangle = \frac{1}{2} \delta_{k, 2p+l}$ .

So we have final result:

$$\sum_k h_k h_{k-2p} = \delta_{0p}$$

**Important sequence** the length of the seq of nonzero  $h_k$  must be even.

### 4. $\sum$ even-indexed = $\sum$ odd-indexed

$$\sum_k h_{2k} = \sum_k h_{2k+1} = \frac{1}{\sqrt{2}}$$

**Proof** Define

$$K_0 = \sum_k h_{2k}, \quad K_1 = \sum_k h_{2k+1}$$

use orthogonality:

$$\sum_k h_k h_{k+2n} = \delta_{0n}$$

Sum both sides over  $n$ :

$$\sum_n \sum_k h_k h_{k+2n} = \sum_n \delta_{0n} = 1$$



Then split the sum:

$$\begin{aligned}
 \sum_n \sum_k h_k h_{k+2n} &= \sum_n \left[ \sum_k h_{2k} h_{2k+2n} + \sum_k h_{2k+1} h_{2k+1+2n} \right] \\
 &= \sum_k \left[ \sum_n h_{2k+2n} \right] h_{2k} + \sum_k \left[ \sum_n h_{2k+2n+1} \right] h_{2k+1} \\
 &= \sum_k \left[ \sum_n h_{2(k+n)} \right] h_{2k} + \sum_k \left[ \sum_n h_{2(k+n)+1} \right] h_{2k+1} \\
 &= K_0 \sum_k h_{2k} + K_1 \sum_k h_{2k+1} \\
 &= K_0^2 + K_1^2 \\
 &= 1
 \end{aligned}$$

and  $K_0 + K_1 = \sqrt{2}$ , then  $K_0 = K_1 = \frac{1}{\sqrt{2}}$

### 6.1.4 Relating the support of $\phi(x)$ to nonzero $h_k$ coeff

**Theorem** If  $\phi(x)$  finite support on  $[N_1, N_2]$ , both  $\mathbb{Z}$ , then  $h_k = 0$  for both  $k > N_2$  and  $k < N_1$ .

In this case,  $h_k$  are said to have compact support in  $[N_1, N_2]$ .

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Is proof required?

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**Proof** the support of the function  $\phi(2x - k)$  must lie inside the interval determined by the ineq:

$$N_1 \leq 2x - k \leq N_2 \implies \frac{1}{2}(N_1 + k) \leq x \leq \frac{1}{2}(N_2 + k)$$

Also, interval must lie inside  $[N_1, N_2]$ :

$$\left[\frac{1}{2}(N_1 + k), \frac{1}{2}(N_2 + k)\right] \subseteq [N_1, N_2] \implies N_1 \leq k \leq N_2$$

**Theorem** If  $\phi(x)$  finite  $[N_1, N_2]$ , then  $\psi(x)$  compact on  $[\frac{1}{2}(N_1 - N_2 + 1), \frac{1}{2}(N_2 - N_1 + 1)]$ .

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Is proof required?

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## Week 12

### 7.1 MRA: general treatment (cont'd)

#### 7.1.1 $\psi(x)$ its support, and vanishing moments

$$\begin{aligned}\int_{\mathbb{R}} \psi(x) dx &= \sum_k g_k \sqrt{2} \int_{\mathbb{R}} \phi(2x - k) dx \\ &= \sum_k (-1)^k h_{1-k} \sqrt{2} \int_{\mathbb{R}} \phi(2x - k) dx\end{aligned}$$

for any  $k \in \mathbb{Z}$ :

$$\begin{aligned}\int_{\mathbb{R}} \phi(2x - k) dx &= \frac{1}{2} \int_{\mathbb{R}} \phi(x) dx \quad (s = 2x - k, ds = 2dx, \text{ etc. } ) \\ &= \frac{1}{2} M\end{aligned}$$

Therefore

$$\begin{aligned}\int_{\mathbb{R}} \psi(x) dx &= \frac{1}{\sqrt{2}} M \sum_k (-1)^k h_{1-k} \\ &= -\frac{1}{\sqrt{2}} M \sum_l (-1)^l h_l \quad (l = 1 - k \Rightarrow k = 1 - l, \text{ etc. } ) \\ &= -\frac{1}{\sqrt{2}} M \left[ \sum_k h_{2k} - \sum_k h_{2k+1} \right] \\ &= 0\end{aligned}$$

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Let us now assume that  $f(x)$  is a polynomial:

$$f(x) = \sum_{k=0}^n c_k x^k$$

over interval  $I$  which contains the domain of support of  $D$  of a wavelet  $\psi$ .

$$\begin{aligned} b_{00} &= \int_I f(x)\psi(x)dx \\ &= \int_D \sum_{k=0}^n c_k x^k \psi(x) dx \\ &= \sum_{k=0}^n c_k \int_D x^k \psi(x) dx \\ &= \sum_{k=0}^n c_k m_k \end{aligned}$$

where

$$m_k = \int_{\mathbb{R}} x^k \psi(x) dx \quad k \geq 0$$

is known as  $k$ th moment of the wavelet function. If  $m_k = 0$  for  $0 \leq k \leq n$  then  $b_{00} = 0$ .

**Defn**  $\psi$  has  $M$  vanishing moments:

$$m_k = \int_{\mathbb{R}} x^k \psi(x) dx = 0, \quad k = 0, 1, 2, \dots, M-1$$

Implications:  $M-1$  vanishing moment, we have an upper bound to  $b_{jk}$

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How much should we know about this?

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### 7.1.2 Vanishing moments and the approximation of functions

**Theorem**  $V_j$  with scaling function  $\phi$ .  $\psi$  has  $M$  vanishing moments. Best approx in  $V_j$  in  $L^2$  sense:

$$f_j = \sum_{k \in \mathbb{Z}} \langle f, \phi_{jk} \rangle \phi_{jk}$$

Then  $L^2$  error has an bound

$$\|f - f_j\|_2 \leq C2^{-jM}$$

## 7.2 MRA and Fourier transform

Two theorems here (not responsible)

### 7.2.1 Fourier transforms and vanishing moments of wavelets

$$\begin{aligned} F^{(n)}(\omega) &= \frac{d^n}{d\omega^n} \left[ \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x) e^{-i\omega x} dx \right] \\ &= (-i)^n \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} x^n f(x) e^{-i\omega x} dx \end{aligned}$$

From this result,

$$\Psi^{(k)}(0) = 0 \quad \text{implies that } m_k = 0$$

Therefore, if

$$\Psi^{(k)} = 0, \quad 0 \leq k \leq M - 1$$

then  $\psi(x)$  has  $M$  vanishing moments.